

Université de Montréal, Canada



ABINIT Workshop 2011

P

- Reminder of the theory
- MLWF
- Code explanation
- Iron pnictides superconductors
- Results
- Magnetic Breakdown

Plan

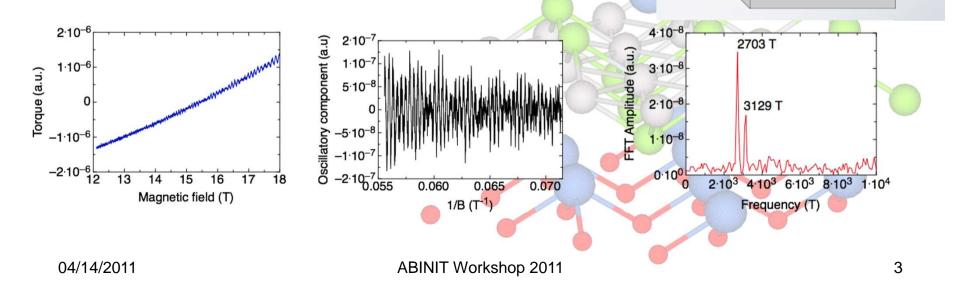
#### Experiment

Cantilever

Sample

Mobile plate Fixe plate

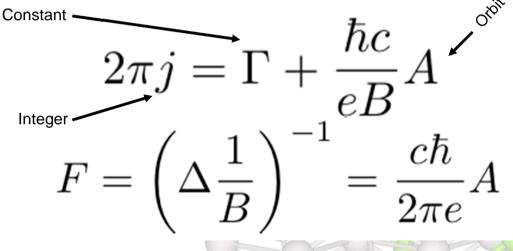
- Sample in a magnetic field
- Torque on the cantilever measures the sample's magnetization
- Oscillations are important!



See Condensed Matter Physics, M.P. Marder, 2000

#### Semi-classical equations

Bohr-Sommerfeld quantization



Ashcroft & Mermin, Solid State Physics (1976)

- Frequencies are a measure of the Fermi surface!
- Only extrema contribute.

### Precision

- Typical ab initio calculation for copper: 6x6x6 k-points
- Too few points to calculate the area of a cross-section precisely
- We need interpolation.
- No need to sample all the Brillouin zone: we only need information on a plane normal to the magnetic field.

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### MLWF

- From Abinit -> Wannier90
- Generalized Wannier function:

$$|u_{n\mathbf{k}}\rangle \rightarrow \sum U_{mn}^{(\mathbf{k})}|u_{m\mathbf{k}}\rangle$$

• With U<sup>(k)</sup> defined such that the spread is minimized<sup>1,2</sup>

$$\sum_{n} \sum_{\mathbf{k}} \langle u_{n\mathbf{k}} | r^2 | u_{n\mathbf{k}} \rangle - | \langle u_{n\mathbf{k}} | \mathbf{r} | u_{n\mathbf{k}} \rangle |^2$$

 The Hamiltonian in this rotated basis (U) is no longer diagonal:

$$\epsilon_{n\mathbf{k}}\delta_{nm} \to H_{nm}(\mathbf{k})$$

1. Marzari N. and Vanderbilt D., PRB **56**, 12847 (1997) 2. Souza I., Marzari N. and Vanderbilt D., PRB **65**, 035109 (2001)

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#### • Fourier transform : $H({f k}) o H({f R})$

 Since the basis is localized, H(R) converges to zero quickly.

MLWF(2)

- Can be interpolated on an arbitrary point in k-space (called  $\mathbf{q}$ ):  $H(\mathbf{R}) \to H(\mathbf{q})$
- This method is used to calculate the band energies (eigenvalues of H(q) on a plane normal to the magnetic field)

# Getting a faster interpolation

К.

$$H(\mathbf{q}) = \sum_{\mathbf{p}} e^{i\mathbf{q}\cdot\mathbf{R}} H(\mathbf{R})$$

- Number of R = number of k-points in the ABINIT calculations
- Size of the Hamiltonian = number of Wannier functions
  - Tuned by rejecting low (and high) energy state
- Careful about reproducing correctly the FS

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## Code: dHvA.py

- Input:
  - H(R) from Wannier90
  - Direction for the magnetic field
  - Option for the Fermi surface slice
- Output:
  - Energy of the plane
  - Possible orbits (open and close)
  - dHvA frequencies (area of closed orbits)
- Going further:
  - Relative intensity of the signal
  - Effective mass

Input H(R)Option for the cut **Build k-points** H(k) and  $\varepsilon(k)$  in a C module Find the orbits Reject the open ones Output Area (frequencies)

#### Slicing the Fermi surface

- Periodic boundary conditions:
  - Finds all the possible orbits
  - Force a specific direction for B
  - Unable to calculate efficiently the area of a specific orbit (too much useless information)
- 'Manual' choice
  - Manually choose the direction of the magnetic field, the size and origin of the plane
  - Cannot find an orbit (need to know its position)
  - Much more accurate (useless information is cut out)

#### Periodic boundary condition: example BCT crystal

- For a magnetic field in the [cosθ 0 sinθ] direction
- $(\mathbf{a}_1 \ \mathbf{a}_2 \ \mathbf{a}_3) = \begin{pmatrix} -a/2 & a/2 & a/2 \\ a/2 & -a/2 & a/2 \\ c/2 & c/2 & -c/2 \end{pmatrix} \qquad (\mathbf{b}_1 \ \mathbf{b}_2 \ \mathbf{b}_3) = \begin{pmatrix} 0 & \tilde{a}/2 & \tilde{a}/2 \\ \tilde{a}/2 & 0 & \tilde{a}/2 \\ \tilde{c}/2 & \tilde{c}/2 & 0 \end{pmatrix}$ • For y direction, take a segment of length  $\tilde{a}$
- For x&z:

 $(L\cos\theta \ 0 \ L\sin\theta) = n_1\mathbf{b}_1 + n_2\mathbf{b}_2 + n_3\mathbf{b}_3$  $\tan\theta = \frac{\tilde{c}}{\tilde{a}}\frac{n_1 + n_3}{n_3 - n_1} = \frac{a}{c}\frac{p}{q}$ 

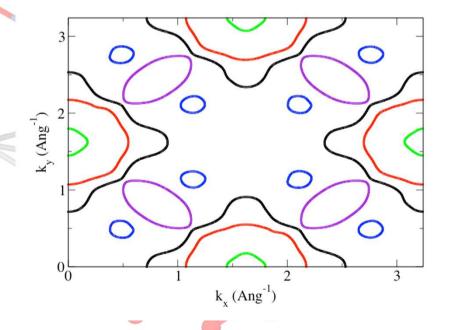
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#### H(q) and diagonalize

- Calculate H(q) for every point on the plane
- Done in a C module: more efficient
- Only interested in a limited set of bands (do not store those not crossing the Fermi level)
- Suppose that each band can be treated separately (disconnected pieces of the FS)

#### Finding the orbit

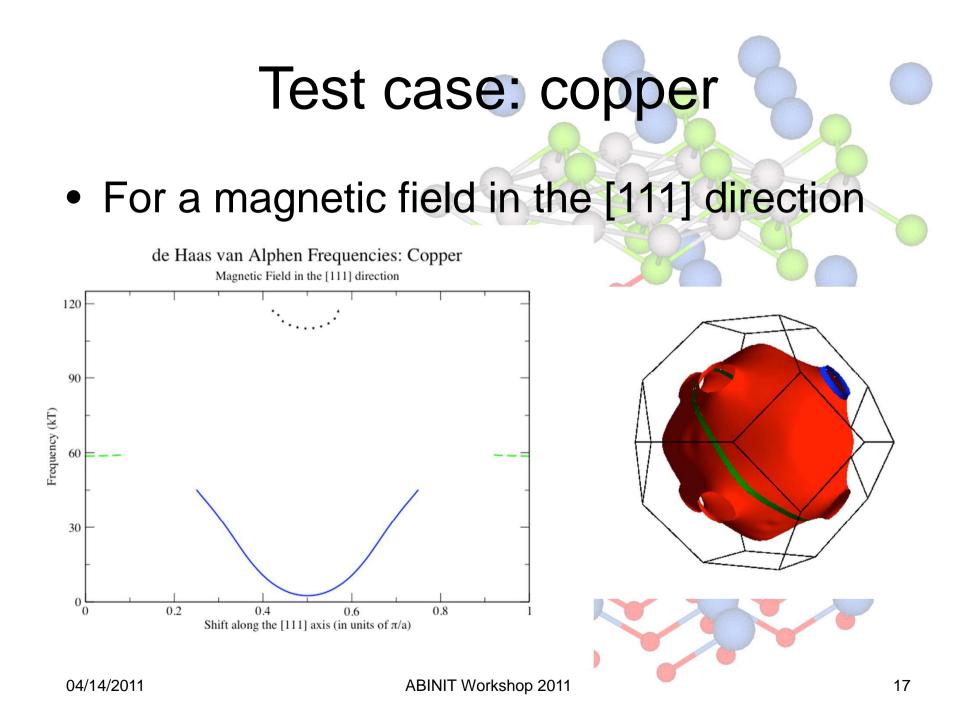
- For a given band, find all the points where the energy changes sign (0 is the Fermi level)
  LaFe<sub>2</sub>P<sub>2</sub>: B=[001]
  k<sub>2</sub>=0.05\*2π/c
- Connect the dots
- Calculate the area
- Converge!



#### Shifting the plane

- Shift the origin of the plane
- Find the extremum: these are the dHvA frequencies
- Intensity (finite difference):  $I \propto \left(\frac{\partial^2 A}{\partial L^2}\right)^{-1/2}$
- Effective cyclotron mass (finite difference):

$$m^* = \frac{\hbar^2}{2\pi} \left(\frac{\partial A}{\partial E_F}\right)$$

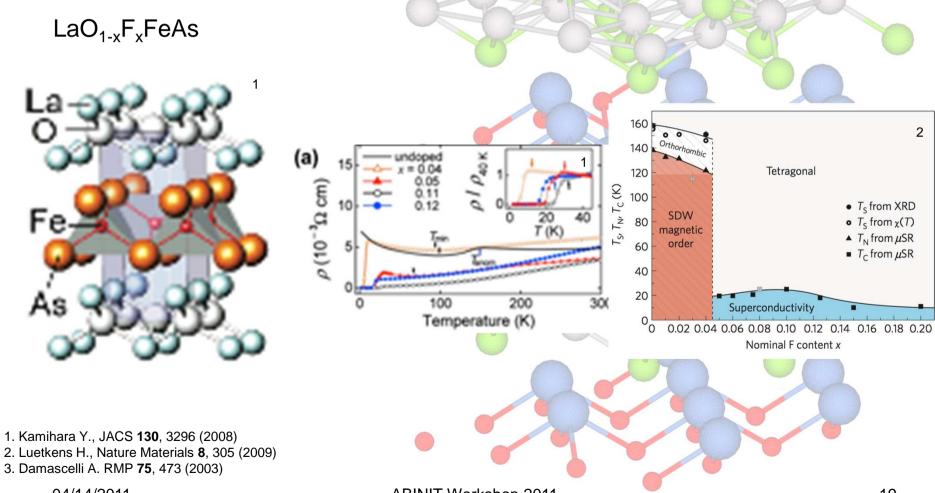


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#### AYNTKAIPSBWTATA

All You Need To Know About Iron Pnictide Superconductor But Were Too Afraid To Ask

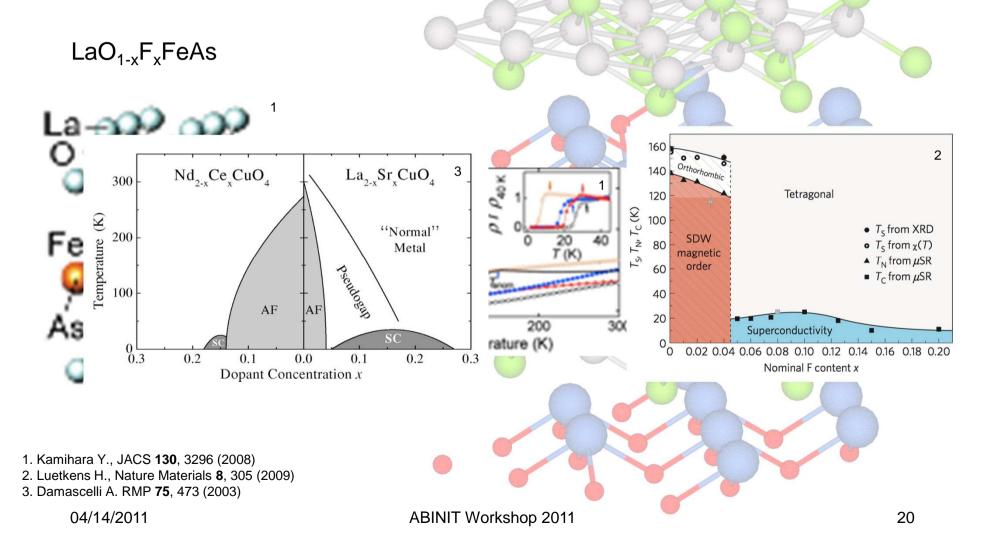


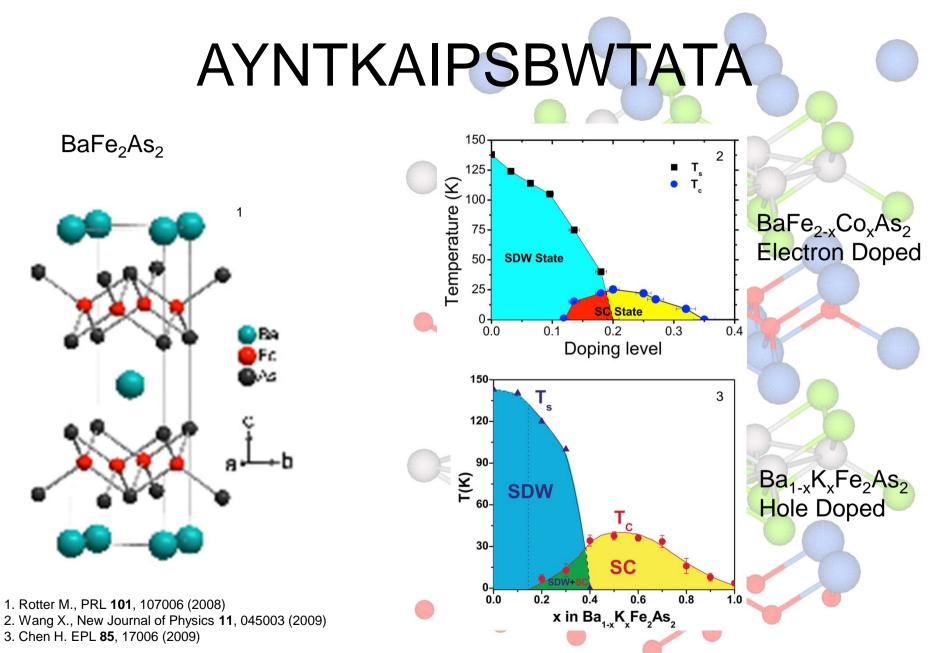
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#### AYNTKAIPSBWTATA

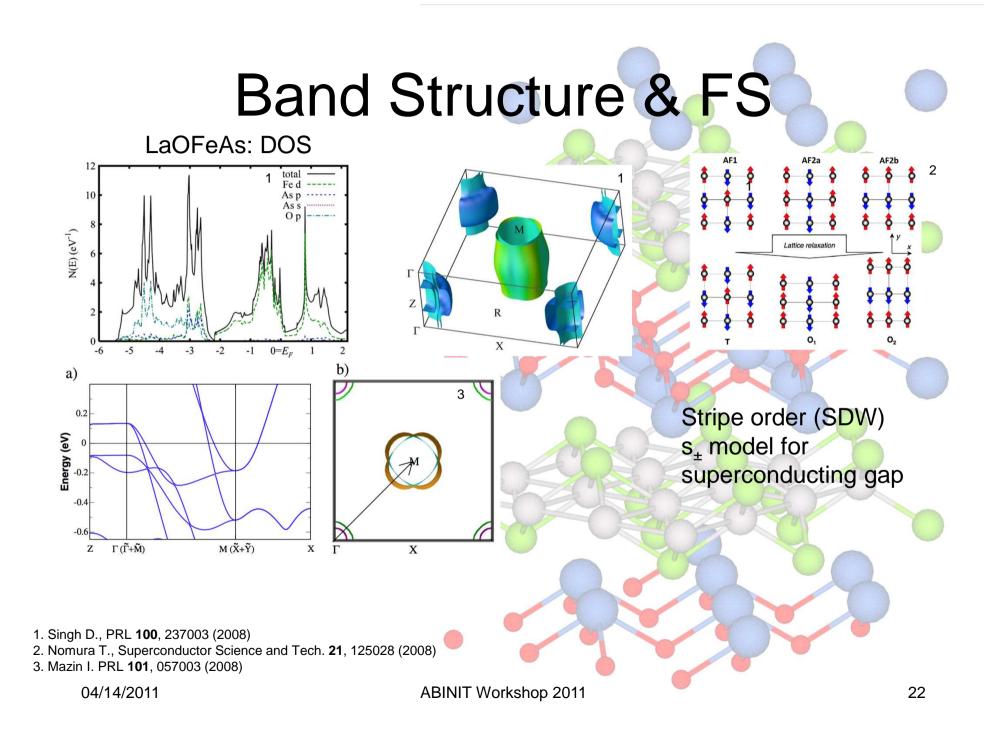
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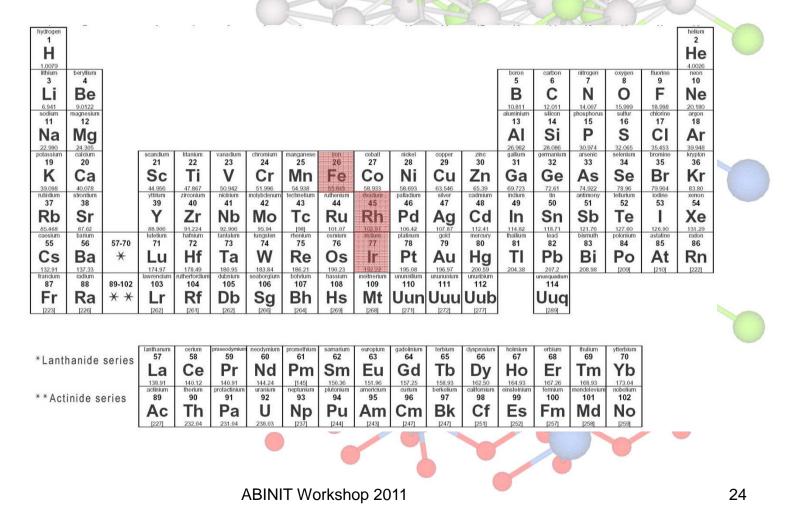


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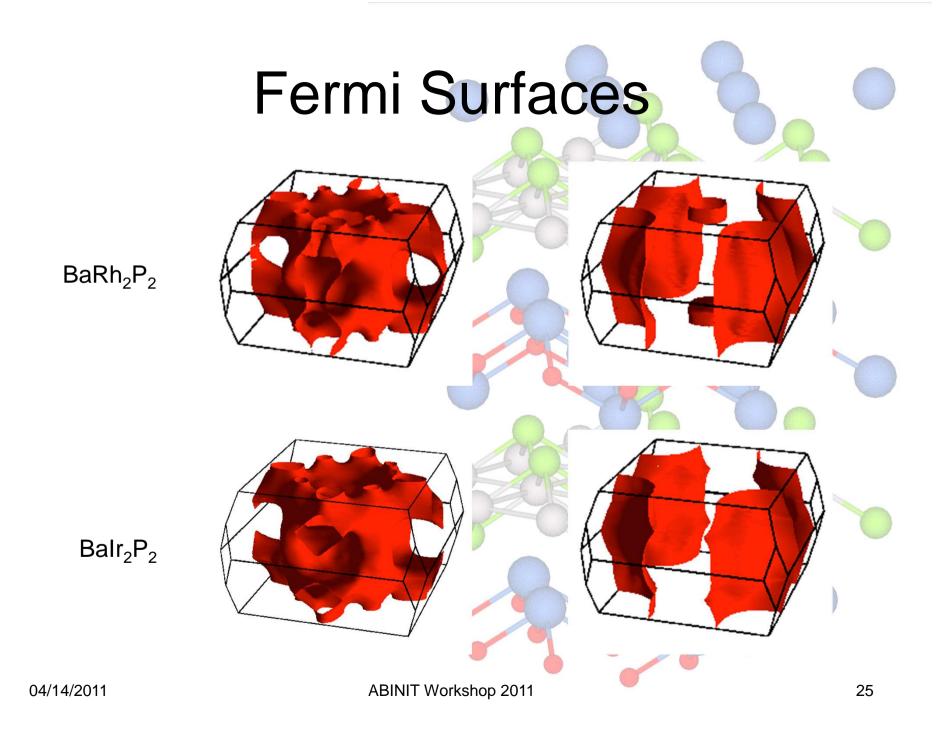
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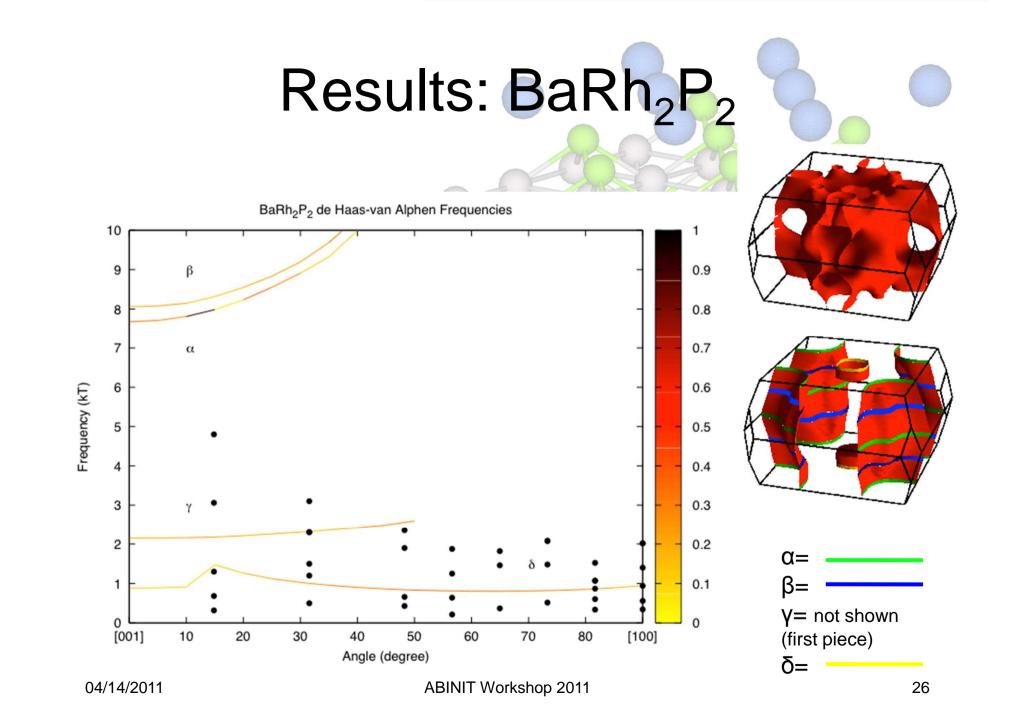
### BaRh<sub>2</sub>P<sub>2</sub> and Balr<sub>2</sub>P<sub>2</sub>

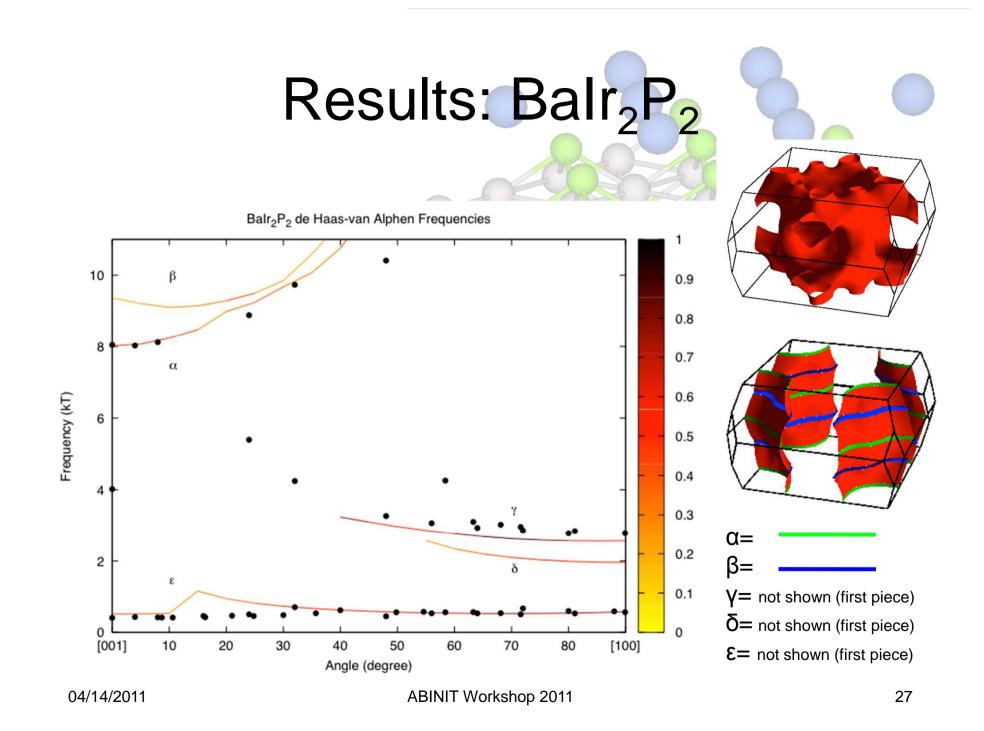
We substitute iron (BaFe<sub>2</sub>P<sub>2</sub>) by rhodium or iridium.



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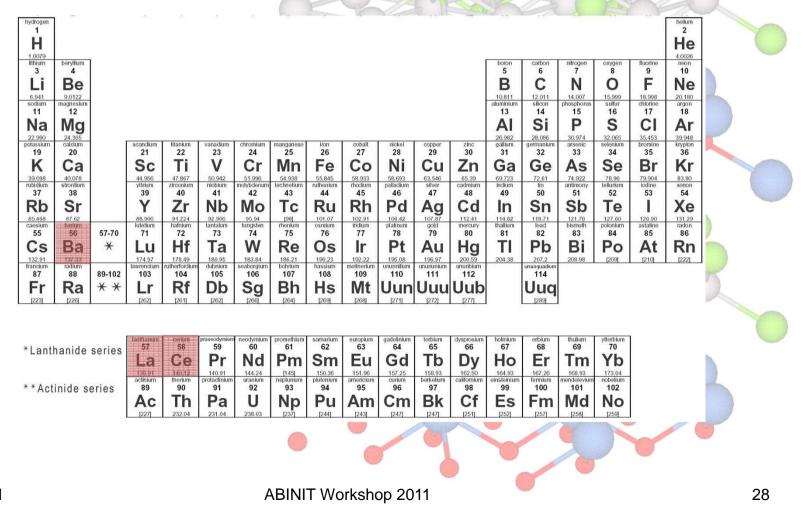


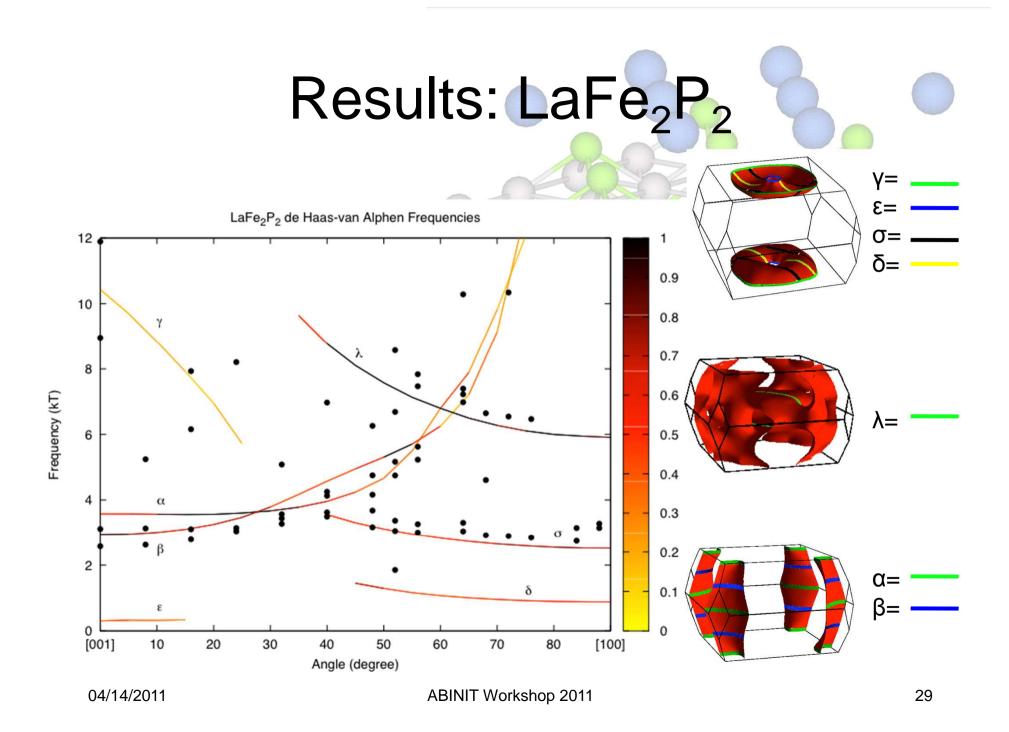


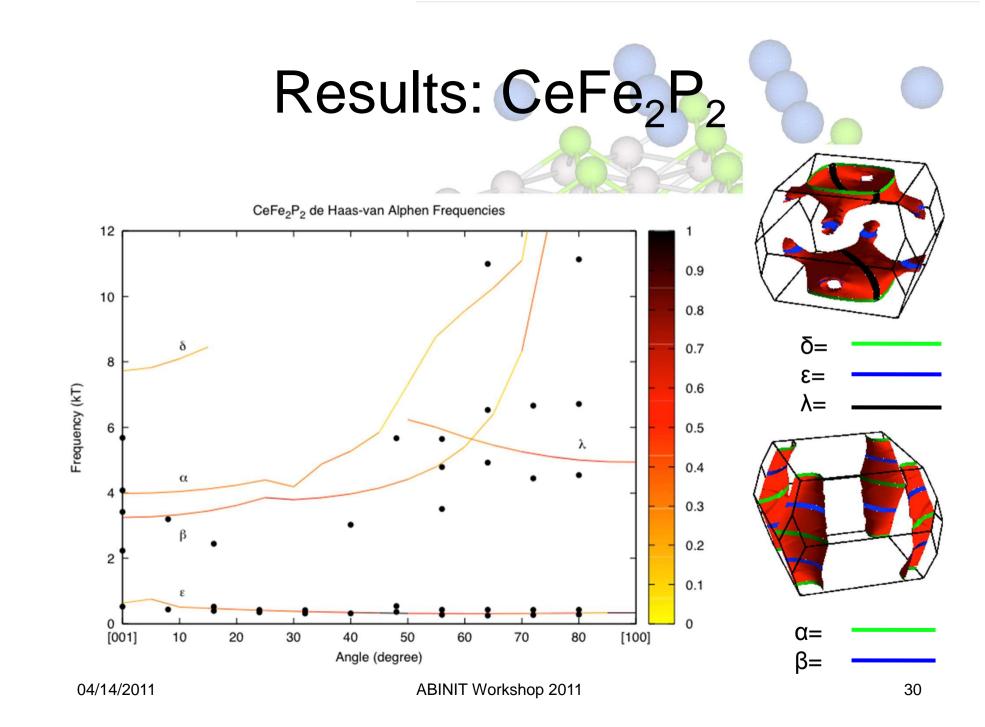


### LaFe<sub>2</sub>P<sub>2</sub> and CeFe<sub>2</sub>P<sub>2</sub>

Changing the rare earth





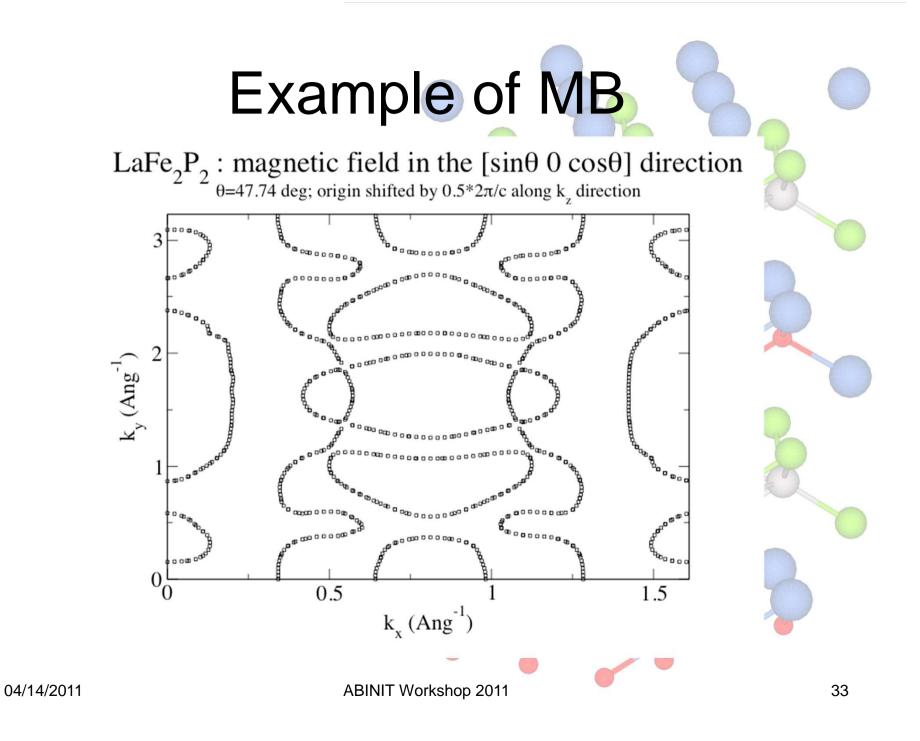


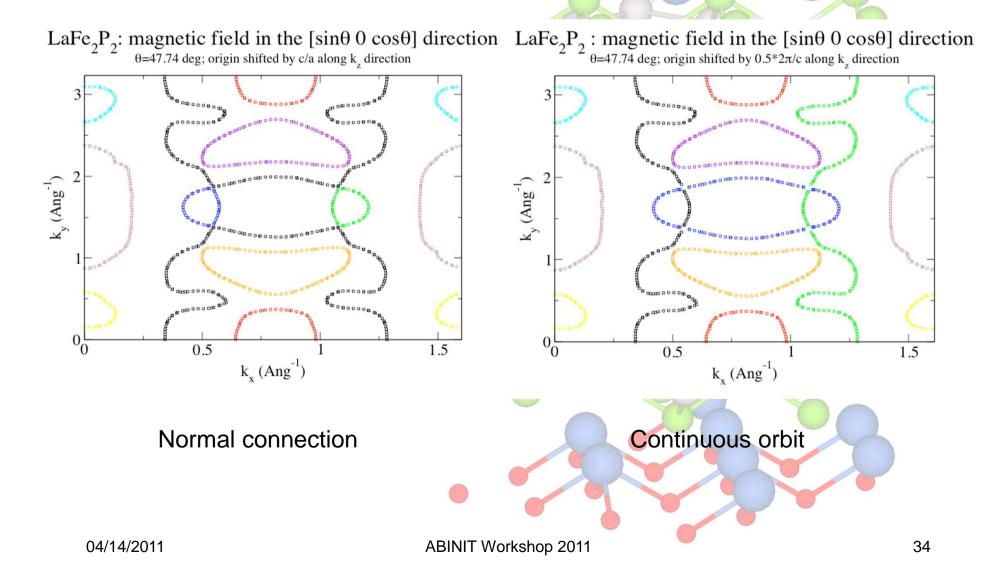
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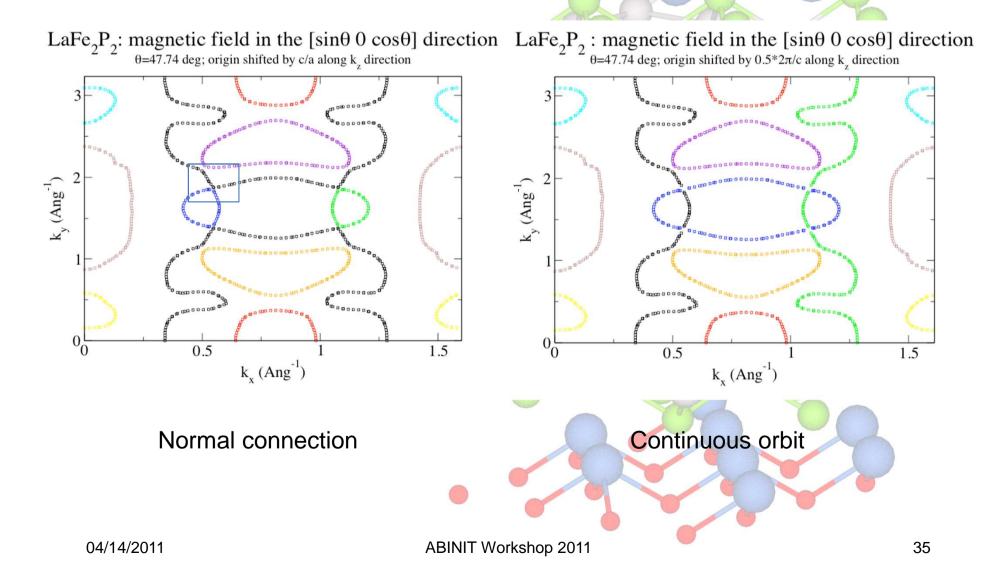
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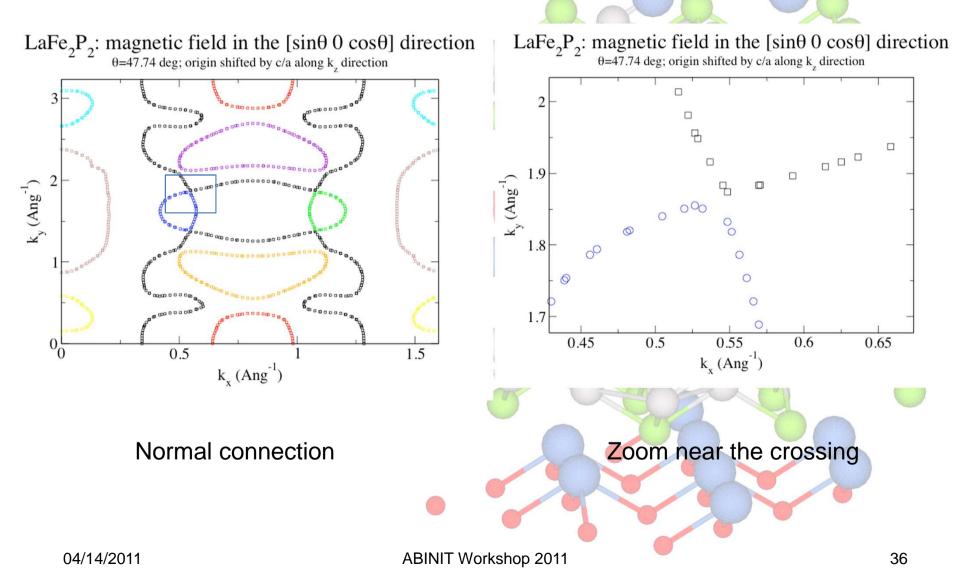
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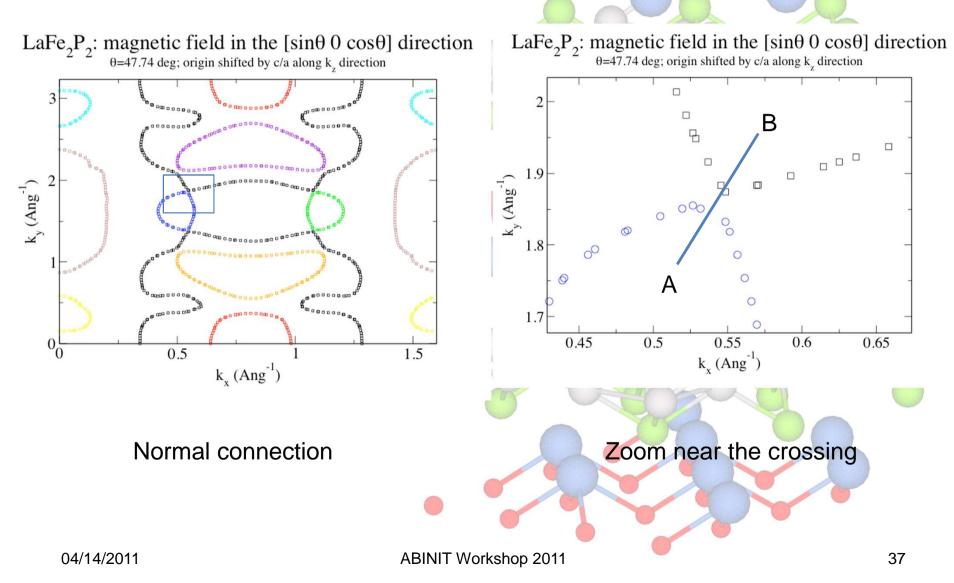
 When 2 pieces of the FS are close, the orbit can change

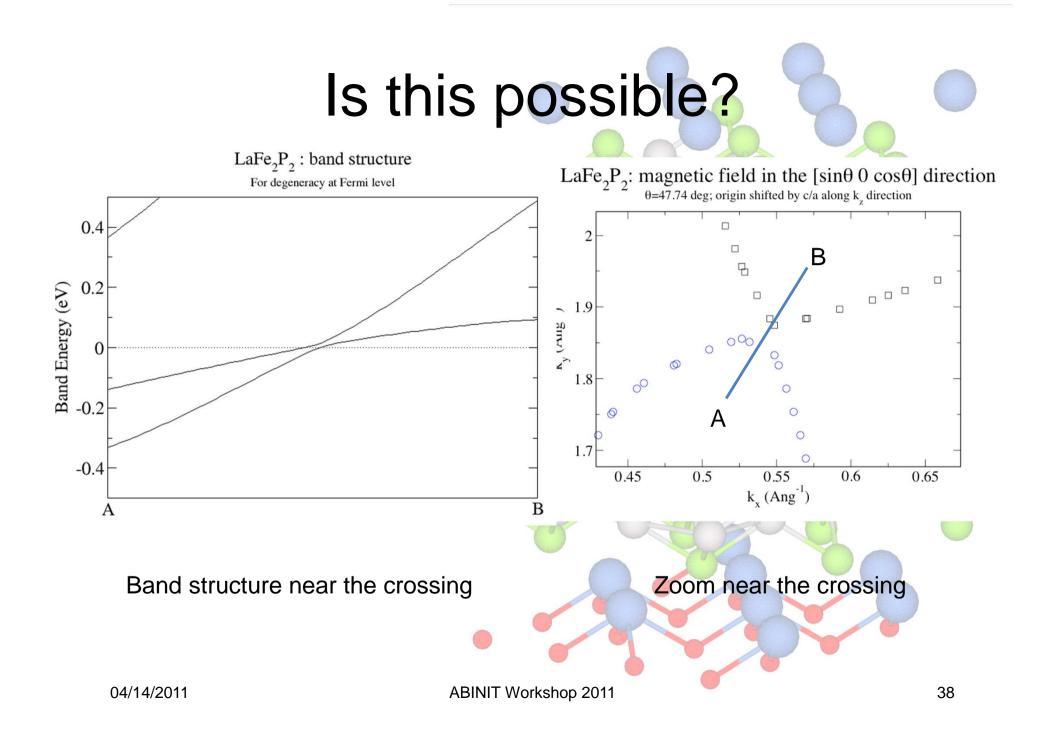












#### Conclusion

- We presented a novel way of calculating dHvA frequencies using maximally localized Wannier function
  - Accurate
  - Low computational cost
- Results for various iron pnictides: LaFe2P2, CeFe2P2, BaRh2P2, Balr2P2
- If 2 (or more) bands are almost degenerate at the Fermi level, the electron orbit is reconstruted (breakdown of the semi-classical equation of motion)

#### Acknowledgement



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