The electron-phonon interaction in ABINIT

## J.-P. Crocombette, X. Gonze, M. Giantomassi, M. Verstraete

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## Why care about el-phon?

■ Phonons are the main scattering mechanism for $\mathrm{T}>0$

- Thermal properties
- Resistance
- Molecular conduction
- Superconductivity


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## Outline

## 1 Basics

## 2 Tutorial

## 3 Novelties

4 Example

5 Conclusions

## Migdal

- Use the Migdal approximation:


■ Separate the explicit coupling term (Frölich type Hamiltonian)

$$
\begin{gathered}
\hat{H}=\hat{H}_{e l}+\hat{H}_{p h}+\hat{H}_{e-p h} \\
\hat{H}_{e-p h}=\sum_{k q}\langle k+q| \nabla_{\alpha} V|k\rangle u_{q \alpha} c_{k+q}^{\dagger} c_{k} \\
\nabla V=\epsilon^{-1} \nabla V_{0} \quad \vec{u}_{q}=\sum_{i} \sqrt{\frac{\hbar}{2 N M \omega_{q i}}} \vec{\epsilon}_{q i}\left(a_{q i}+a_{q i}^{\dagger}\right)
\end{gathered}
$$

■ The self-energy for the phonons is in the LR screening - The self-energy for the electrons


■ Eliashberg function (weighted DOS)


## Eliashberg

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■ The self-energy for the electrons

$$
\Sigma_{e p}=T \int_{F S} \int_{\Omega} \frac{\alpha^{2} F\left(k, k^{\prime}, \Omega\right)}{N(0)}\left(\frac{2 \Omega}{\omega^{2}+\Omega^{2}}\right) G
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■ Eliashberg function (weighted DOS)

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\begin{gathered}
\alpha^{2} F\left(k, k^{\prime}, \Omega\right)=N(0) \sum_{j}\left|g_{k, k^{\prime}}^{j}\right|^{2} \delta\left(\omega_{q, j}-\Omega\right) \\
g_{k, k^{\prime}}^{j}=\frac{1}{\sqrt{2 M \omega_{q, j}}} \vec{\epsilon}_{q, j} \cdot\langle k| \nabla V\left|k^{\prime}\right\rangle
\end{gathered}
$$

## EP quantities

■ EP coupling strength (anisotropic)

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\lambda\left(k, k^{\prime}, \omega\right)=\int_{0}^{\infty} d \Omega \frac{2 \Omega}{\omega^{2}+\Omega^{2}} \alpha^{2} F\left(k, k^{\prime}, \Omega\right)
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\gamma_{q, j}=2 \pi \omega_{q, j} \int_{F S}\left|g_{k+q, k}^{q j}\right|^{2}
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## Superconductivity

$\square$ McMillan equation is popular

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T_{c}=\frac{\omega_{\log }}{1.2} \exp \left(\frac{-1.04(1+\lambda)}{\lambda-\mu^{\star}(1+0.62 \lambda)}\right)
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■ In linear response the coupling is byproduct of Sternheimer equation $\rightarrow H^{(1)}\left|\psi_{0}\right\rangle$

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## Demonstration with the tutorial

■ telphon_1 calculate GS and all the $3^{*} \mathrm{~N}_{\text {atom }}$ phonons

- telphon_2 merge the DDB files (mrgddb)
- telphon_3 merge the GKK files (mrggkk)
- telphon 4 run anaddb


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■ On selected BS qpoints interpolate BS and linewidths

- On dense grid calculate $\alpha^{2} F(\Omega)$
$\square$ Calculate moments of $\alpha^{2} F(\Omega)$
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## Goodies from Matteo Giantomassi

■ FS output

- Nesting factor
- Interpolation checks

$\square \longrightarrow$ see next talk, with examples!


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Higher moments of a2F

■ Standard $\lambda$ is special casea ( $\mathrm{n}=0$ ) of

$$
\begin{equation*}
\lambda<\omega^{n}>=2 \int_{0}^{\infty} d \Omega\left[\alpha^{2} F(\Omega)\right] \Omega^{n} \tag{1}
\end{equation*}
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- Added calculation of $\lambda<\omega^{n}>$ for $n=2,3,4,5$
- Used to estimate the temperature relaxation rate of hot electrons
from Allen PRL 591460 (1987)

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\begin{equation*}
\gamma_{T}=\frac{3 \hbar \lambda<\omega^{2}>}{\pi k_{B} T_{e}}\left(1-\frac{\hbar^{2} \lambda<\omega^{4}>}{12 \lambda<\omega^{2}>k_{B}^{2} T_{e} T_{L}}+\ldots\right) \tag{2}
\end{equation*}
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■ Interpolating phonons and elphon matrices separately

- Basis at random qpoint could be different
- Standard diagonalization routines give monotonically increasing order for $\omega$ and $\gamma_{q j}$


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## Mode separation II

Straight diag of dynamical matrix and $\gamma$ matrix


## Mode separation III

■ Which linewidth belongs to which phonon mode?

- Interpolate perturbations before scalar product w/ $\overrightarrow{q i}$

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## Phonon resistivity contribution I

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- Momentum relaxation due to scattering off phonons



## Phonon resistivity contribution II

Transport spectral function $\alpha_{t r} F=\alpha_{\text {out }}^{2} F-\alpha_{\text {in }}^{2} F$

$$
\begin{aligned}
& \alpha_{\text {out }}^{2} F(\omega)=\frac{1}{N(0)\left\langle v_{x}^{2}\right\rangle} \sum_{\nu} \sum_{k k^{\prime} j^{\prime}}\left|g_{q \nu}^{k k^{\prime} j^{\prime}}\right|^{2} v_{x}(k) v_{x}(k) \delta\left(\epsilon_{k j}\right) \delta\left(\epsilon_{k^{\prime} j^{\prime}}\right) \delta\left(\omega-\omega_{q \nu}\right) \\
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$$
\rho(T)=\frac{\pi \Omega_{\text {cell }} k_{B} T}{N(0)\left\langle v_{x}^{2}\right\rangle} \int_{0}^{\infty} \frac{d \omega}{\omega} \frac{x^{2}}{\sinh ^{2}(x)} \alpha_{t r} F(\omega) \quad x=\frac{\omega}{k_{B} T}
$$

## Limitations/problems

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$$
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$\square$ Phase disappears in $\left|g_{k n, k^{\prime} n^{\prime}}^{q j}\right|^{2}$

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## FCC lead

- Compare Eliashberg function with litterature
- Spin-orbit coupling is essential




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## State of the art

1 Fully functional el-phon code. Several papers published (i.e. not just by me)

2 Potential for many extensions: superconductivity, (anisotropic) transport
3 Input for Raman spectra? Higher order or resonant Raman
4 And beyond...
5. Thank you for your attention!

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## Beyond BO and Eliashberg

Existing extensions of the standard model
$\square$ Gross or VanLeeuwen formalisms for el and nucleus density matrix
> - Include e-e and e-p interactions together diagrammatically

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